

# ON THE FUNDAMENTAL INVARIANT OF THE HECKE ALGEBRA $H_n(q)$

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The fundamental invariant of the Hecke algebra  $H_n(q)$  is the  $q$ -deformed class-sum of transpositions of the symmetric group  $S_n$ . Irreducible representations of  $H_n(q)$ , for generic  $q$ , are shown to be completely characterized by the corresponding eigenvalues of  $C_n$  alone. For  $S_n$  more and more invariants are necessary as  $n$  increases. It is pointed out that the  $q$ -deformed classical quadratic Casimir of  $SU(N)$  plays an analogous role. It is indicated why and how this should be a general phenomenon associated with  $q$ -deformation of classical algebras. Apart from this remarkable conceptual aspect  $C_n$  can provide powerful and elegant techniques for computations. This is illustrated by using the sequence  $C_2, C_3, \dots, C_n$  to compute the characters of  $H_n(q)$ .

This talk will be based on [1] and [2]. Much more complete discussions and references to other authors can be found there. More recent developments can be found in [3].

Let me start by recapitulating certain facts concerning the invariants of the classical symmetric group  $S_n$ . The single cycle class-sums  $\{[p]_n; p = 2, 3, \dots, n\}$  belong to the centre. Here  $[2]_n$  is the sum of transpositions,  $[3]_n$  is that of circular permutation of triplets (each term being a product of transpositions) and so on. Their eigenvalues characterize irreducible representations (irreps.) of  $S_n$  corresponding to different standard Young tableaux with  $n$  boxes.

**Definition.** *Content of the box in the  $i$ -th row and  $j$ -th column of the Young tableau is equal to  $(j - i)$ .*

The symmetric power sums of these contents gives the eigenvalues  $\lambda_{[p]_n}^\Gamma$  of  $[p]_n$  for  $\Gamma$ . Thus

$$\lambda_{[2]_n}^\Gamma = \sum_{(i,j) \in \Gamma} (j - i) \quad (1)$$

$$\lambda_{[3]_n}^\Gamma = \sum_{(i,j) \in \Gamma} (j - i)^2 - \frac{1}{2}n(n - 1) \quad (2)$$

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and so on [1].

For  $n \geq 6$  the eigenvalues of  $[2]_n$  show degeneracy. As  $n$  increases higher and higher  $[p]_n$ 's are needed to uniquely characterize each irreducible representation.

*The situation changes dramatically as  $S_n$  is  $q$ -deformed to  $H_n(q)$ . The eigenvalues of  $C_n$ , the  $q$ -deformed  $[2]_n$ , alone suffice to characterize the irreducible representations for arbitrary  $n$ . (Throughout only real, positive i.e. generic  $q$  is considered.)*

I will now show how this becomes possible. The generators of the  $H_n(q)$  satisfy

$$\begin{aligned} g_i^2 &= (q-1)g_i + q & i &= 1, 2, \dots, n-1 \\ g_i g_{i+1} g_i &= g_{i+1} g_i g_{i+1} & i &= 1, 2, \dots, n-2 \\ g_i g_j &= g_j g_i & \text{if } |i-j| &\geq 2 \end{aligned} \quad (3)$$

For  $q = 1$  one gets  $S_n$ . The fundamental invariant is

$$\begin{aligned} C_n &= g_1 + g_2 + \dots + g_{n-1} + \frac{1}{q}(g_1 g_2 g_1 + g_2 g_3 g_2 + \dots + g_{n-2} g_{n-1} g_{n-2}) \\ &+ \frac{1}{q^2}(g_1 g_2 g_3 g_2 g_1 + g_2 g_3 g_4 g_3 g_2 + \dots + g_{n-3} g_{n-2} g_{n-1} g_{n-2} g_{n-3}) \\ &+ \dots \\ &+ \frac{1}{q^{n-2}} g_1 g_2 \dots g_{n-2} g_{n-1} g_{n-2} \dots g_2 g_1 \end{aligned} \quad (4)$$

For  $q = 1$ ,  $g_1 g_2 g_1 = (13)$  and so on and one gets back  $[2]_n$ . The eigenvalue of the fundamental invariant for the  $Y$ -tableau  $\Gamma$  can be shown [1] to be the following  $q$ -deformation of (1),

$$\Lambda_n^\Gamma = q \sum_{(i,j) \in \Gamma} \frac{q^{j-i} - 1}{q - 1} = q \sum_{(i,j) \in \Gamma} [j - i]_q. \quad (5)$$

**Definition.**  $q$ -content of the box in the  $i$ -th row and  $j$ -th column of the Young tableau is equal to  $q [j - i]_q$ .

Hence  $\Lambda_n^\Gamma$  is the sum of the  $q$ -contents of the boxes of  $\Gamma$ .

Consider, for  $n = 6$ , the irreducible representations  $[4, 1, 1]$  and  $[3, 3]$ . The box contents are

0	1	2	3
-1			
-2			

$[4, 1, 1]$

0	1	2
-1	0	1

$[3, 3]$

For  $S_6$ ,

$$\Lambda_6^{[4, 1, 1]} = \Lambda_6^{[3, 3]} = 3 \quad (6)$$

For  $H_6(q)$ ,

$$\begin{aligned} \Lambda_6^{[4, 1, 1]}(q) &= q^3 + 2q^2 + 3q - 2 - \frac{1}{q} \\ \Lambda_6^{[3, 3]}(q) &= q^2 + 3q - 1 \end{aligned} \quad (7)$$

Thus the degeneracy is lifted as  $q$  moves away from unity. This is the simplest non-trivial example. For the general case one notes:

- (i) The  $q$ -contents are constant for boxes on the same diagonal of  $\Gamma$ .
- (ii) Developping the  $q$ -brackets and regrouping terms

$$\Lambda_n^\Gamma = \sum_{k>0} q^k \pi_k^\Gamma - \sum_{k<0} q^{k+1} \nu_k^\Gamma \quad (8)$$

where

$$\begin{aligned} \pi_k^\Gamma &= \sum_{l \geq k > 0} (\text{number of boxes with content } l) \\ \nu_k^\Gamma &= \sum_{l \leq k < 0} (\text{number of boxes with content } l) \end{aligned}$$

From (i) and (ii) it is not difficult to show [1] that  $\Lambda_n^\Gamma$  completely determines  $\Gamma$  and hence the irrep.

Set  $q = e^\delta$  ( $\neq 1$ ) and let

$$\tilde{C}_n \equiv \frac{q-1}{q} C_n \quad (9)$$

then [1],

$$\Lambda_{\tilde{C}_n}^\Gamma = \delta \lambda_{[2]_n}^\Gamma + \frac{\delta^2}{2} (\lambda_{[3]_n}^\Gamma + \frac{1}{2} n(n-1)) + \dots \quad (10)$$

The eigenvalues of all  $[p]_n$  ( $p = 2, \dots, n$ ) in the coefficients of the above series. This is another way of exhibiting that  $C_n$  by itself contains information equivalent to that supplied by all  $[p]_n$  for  $S_n$ .

Projection operators for irreps. can be constructed in terms of  $C_n$  in a straightforward way since there is no degeneracy. When the limit  $q \rightarrow 1$  is taken correctly higher class-sums of  $S_n$  appear automatically as necessary to project out the corresponding irrep. of  $S_n$ . This is discussed in detail in [1]. Further interesting uses of projection operators can be found in [3].

In [1] a direct relation was given between the eigenvalues of  $C_n$  and those of the Casimir of  $SU_q(N)$  ( $q$ -deformation of the Casimir quadratic in the Cartan-Weyl generators of  $SU(N)$ ) for an irrep. corresponding to a  $Y$ -diagram  $\Gamma$  with

$n$  boxes (and at most  $N - 1$  rows). It was shown that this Casimir  $C_2$  can be so redefined (denoted then by  $\tilde{C}_2$ ) that the eigenvalue for  $\Gamma$  is just

$$\Lambda_{\tilde{C}_2}^\Gamma = \sum_{k=1}^{N-1} q^{2(l_k - k)} \quad (11)$$

where  $l_k$  is the number of boxes in the  $k$ -th row. This was derived using the Gelfand-Zetlin basis [1]. But (11) is, of course, independent of the choice of such a basis. Since

$$l_k \geq l_{k+1}, \quad (l_k - k) > (l_{k+1} - k - 1) \quad (12)$$

Hence the indices of  $q$  in (11) are strictly monotonically decreasing.

Thus even for a reducible representation arising in a certain context (say some model) if one obtains the matrix of  $\tilde{C}_2$  from some source and diagonalizes it the coefficient of each block of unit matrix must be of the form

$$\sum_{k=1}^{N-1} q^{2L_k} \quad (L_k > L_{k+1}) \quad (13)$$

Now setting

$$l_k = L_k + k, \quad (k = 1, \dots, N - 1) \quad (14)$$

The  $Y$ -diagram  $\Gamma$  is completely determined. *Thus the eigenvalue of a suitably  $q$ -deformed quadratic Casimir completely characterizes an irrep. of  $SU_q(N)$ .* (For  $q = 1$  or  $SU(N)$  one needs, in general, all the invariants upto order  $N$ .)

Setting  $q = e^\delta$ ,

$$\sum_k q^{2L_k} = 1 + (2\delta) \left( \sum_k L_k \right) + \frac{1}{2!} (2\delta)^2 \left( \sum_k L_k^2 \right) + \dots \quad (15)$$

One can compare (15) with (10).

The coefficients of higher powers of  $\delta$  contain informations equivalent to those of higher order Casimirs of  $SU(N)$ .

I present now, without derivation, the relation between  $C_n$  and  $\tilde{C}_2$  eigenvalues for a  $\Gamma$  with  $n$  boxes [1],

$$\left( \frac{q^2 - 1}{q^2} \right)^2 \Lambda_{C_n(q^2)}^\Gamma + \frac{q^2 - 1}{q^2} n = \Lambda_{\tilde{C}_2}^\Gamma + \frac{q^{2(-N+1)} - 1}{q^2 - 1} \quad (16)$$

for  $(\sum l_k = n \text{ in (11)})$ .

The Hecke and  $q$ -deformed unitary algebras are well-known to be closely related. But the aspect presented here is more general. *Thus the  $q$ -deformed quadratic Casimirs of the other Lie algebras should play analogous roles.* Our investigation is not complete. Here only  $SU_q(N)$  has been studied. However in an accompanying talk [4] the foregoing statement is confirmed for  $SO_q(5)$ . The discussion at the end of [4] gives an idea of the richness of content of the  $q$ -deformed Casimirs.

Apart from such remarkable conceptual aspects,  $C_n$  or, even better, the sequence  $(C_2, C_3, \dots, C_n)$  nested in  $H_n(q)$  can furnish powerful techniques for various goals. As an example, I will indicate below how they can be used to compute characters. A detailed study can be found in [2]. (Another interesting aspect has been studied in [3]).

For the sequence  $H_2(q) \subset H_3(q) \subset \dots \subset H_n(q)$  one defines the Murphy operators

$$L_2 = C_2, \quad L_3 = C_3 - C_2, \dots, \quad L_n = C_n - C_{n-1} \quad (17)$$

One obtains

$$L_p = \sum_{i=1}^{p-1} q^{1-p+i} (g_i g_{i+1} \dots g_{p-1} \dots g_{i+1} g_i) \quad (18)$$

and

$$L_{p+1} = \frac{1}{q} g_p L_p g_p + g_p \quad (19)$$

Basis vectors of an irrep. can be specified by sequences of  $Y$ -diagrams (indicating how successive boxes are added)

$$\Gamma_2 \subset \Gamma_3 \subset \dots \subset \Gamma_n \quad (20)$$

The eigenvalue of  $L_i$  can be shown to be [2]

$$\{\Gamma_i \setminus \Gamma_{i-1}\}_q \equiv q[k_i - p_i]_q \quad (21)$$

where  $\Gamma_i$  is obtained by adding the box  $(k_i, p_i)$  to  $\Gamma_{i-1}$ . The eigenvalue is the  $q$ -content of the last box added. Also

$$tr(L_i)_{\Gamma_n} = \sum_{\Gamma_{n-1} \subset \Gamma_n} tr(L_i)_{\Gamma_{n-1}} \quad (i = 2, 3, \dots, n-1.) \quad (22)$$

and

$$tr(L_n)_{\Gamma_n} = \sum_{\Gamma_{n-1} \subset \Gamma_n} |\Gamma_{n-1}| \{\Gamma_n \setminus \Gamma_{n-1}\}_q. \quad (23)$$

where

$$|\Gamma_{n-1}| = \sum_{\Gamma_{n-2} \subset \Gamma_{n-1}} |\Gamma_{n-2}| = \dim \Gamma_{n-1} \quad (24)$$

For what follows we will need traces of products of *non-consecutive* Murphy operators only. For such products with

$$\alpha_{i+1} \geq \alpha_i + 2$$

$$tr \left( \prod_{i=1}^{\ell} L_{\alpha_i} \right)_{\Gamma_n} = \sum_{\Gamma_{n-1} \subset \Gamma_n} tr \left( \prod_{i=1}^{\ell} L_{\alpha_i} \right)_{\Gamma_{n-1}} \quad (\text{for } \alpha_l < n) \quad (25)$$

$$tr \left( \prod_{i=1}^{\ell} L_{\alpha_i} \right)_{\Gamma_n} = \sum_{\Gamma_{n-1} \subset \Gamma_n} \{ \Gamma_n \setminus \Gamma_{n-1} \}_q tr \left( \prod_{i=1}^{\ell-1} L_{\alpha_i} \right)_{\Gamma_{n-1}} \quad (\text{for } \alpha_l = n) \quad (26)$$

The recursion relations (22) to (26) yield easily the traces of the  $L$ 's and their non-consecutive products [2]. A symbolic program is easy to set up. Taking the trace of each side of (18) and inverting the relation one obtains [2].

$$tr(g_1 g_2 \cdots g_{k-1}) = \left( \frac{q}{q-1} \right)^{k-2} \sum_{i=0}^{k-2} (-1)^i \binom{k-1}{i} tr(L_{k-i}) \quad (27)$$

Similarly, after multiplying both sides of (18) by  $L_m$  (non-consecutive),

$$tr((g_1 g_2 \cdots g_{k-1}) L_m) = \left( \frac{q}{q-1} \right)^{k-2} \sum_{i=0}^{k-2} (-1)^i \binom{k-1}{i} tr(L_{k-i} L_m). \quad (28)$$

Continuing step-wise, with suitable choices of  $m$  at each step, it can be shown [2] that one finally obtains in terms of known traces of the type (25) and (26) traces of the form

$$tr \left( (g_1 g_2 \cdots g_{m_1-1}) (g_{m_1+1} \cdots g_{m_2-1}) \cdots (g_{m_j+1} \cdots g_p) (g_k g_{k+1} \cdots g_{p+r} \cdots g_{k+1} g_k) \right) \quad (29)$$

Here the last factor comes from the last non-consecutive  $L$  ( $L_{p+r-1}$ ). This, in general, has overlapping indices with the preceding factors. At each previous step such overlaps has been assumed to be *reduced* (reexpressed as sums of traces of ordered products the  $g_i$ 's in ascending order of  $i$ ) so that they have no overlap but, possibly, *cuts* (at  $i = m_1, m_2, \dots, m_j$ , say). This reduction procedure, to be applied again to (29), will be briefly described below. Let us

however first are exhibit the simplest results, illustrating general properties. One obtains [2]

$$\text{tr}(g_i) = \text{tr}(L_2), \quad (i = 1, \dots, n-1) \quad (30)$$

$$\text{tr}(g_i g_{i+1}) = \left(\frac{q}{q-1}\right) \left(\text{tr}(L_3) - 2\text{tr}(L_2)\right), \quad (i = 1, \dots, n-2) \quad (31)$$

$$\text{tr}(g_i g_{i+2}) = \frac{1}{q-1} \left( -2q \text{tr}(L_2) + (q+1)^2 \text{tr}(L_3) - (1+q^2) \text{tr}(L_4) + (q-1) \text{tr}(L_2 L_4) \right) \quad (32)$$

The equalities of the traces in each example illustrate a fundamental lemma [2]: *The trace of product of any number of disjoint sequences, in any irrep., depends only on the lengths of the component connected sequences.*

Note that in (32),  $(g_i g_{i+2})$  being disjoint (i.e. a cut at  $i+1$ ) a non-consecutive product  $L_2 L_4$  appears on the right. *Note that the  $L$ 's on the right do not depend on  $n$  (of  $H_n(q)$ ).* This is also a general feature.

When there is no overlap and at most one cut (29) can be reduced relatively easily [2]. Thus

$$\begin{aligned} V_k &\equiv \text{tr} \left( (g_1 \dots g_{k-1}) (g_{k+1} \dots g_p) (g_{p+1} \dots g_{p+r} \dots g_{p+1}) \right) \\ &= \sum_{\ell=0}^{r-1} \binom{r-1}{\ell} q^\ell (q-1)^{r-\ell-1} \text{tr} \left( (g_1 \dots g_{k-1}) (g_{k+1} \dots g_{p+r-\ell}) \right) \end{aligned} \quad (33)$$

(For  $V_0$  with  $k=0$  the first factor is defined to be unity). When there is an overlap we introduce *f-expansions* defined below. (For a full account see Appendix [2]). From,

$$g_i^2 = (q-1)g_i + q$$

one deduces

$$g_i^p = f_p g_i + q f_{p-1} \quad (34)$$

where

$$f_p = \frac{q^p - (-1)^p}{q+1}$$

It can be shown that (for overlap =  $p-l+1$ )

$$\text{tr} \left( (g_1 \dots g_p) (g_l \dots g_{p+r} \dots g_l) \right) = (q-1) \sum_{r=1}^{p-l+1} q^k f_{2(p-l+1-k)+1} V_k + f_{2(p-l+1)+1} V_0 \quad (35)$$

*Thus the  $f$ -coefficients are determined only by the length of the overlap.* The general case with overlap and multiple cuts is treated in App. [2].

Tables of characters (polynomials in  $q$ ) are given in [2]. Here let me just summarize the main steps:

- (1) Traces of (non-consecutive) products of Murphy operators.
- (2) Traces of products of  $g$ 's with cuts and overlap in terms of (1).
- (3) Reduction removing overlaps ( $f$ -expansions).

## References

1. J. Katriel, B. Abdesselam and A. Chakrabarti, The fundamental invariant of the Hecke algebra  $H_n(q)$  characterizes the representations of  $H_n(q)$ ,  $S_n$ ,  $SU(N)_q$  and  $SU(N)$ . (q-alg/9501021)(to published in Jour. Math. Phys).
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3. T. Brzezinski and J. Katriel, representation theoretic derivation of the Temperley-Lieb-Martin algebras (hep-th/9507128).
4. A. Chakrabarti, talk presented in the Nakai workshop (Tianjin, 1995).